

# Exciton g-factors of van der Waals heterostructures from first-principles calculations

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## ABSTRACT

We present how exciton g-factors and their sign can be determined by converged first principles calculations. We apply the method to monolayer excitons in transition metal dichalcogenides (TMDs) and to interlayer excitons in MoSe<sub>2</sub>/WSe<sub>2</sub> heterobilayers. The precision of the method allows to assign measured g-factors of optical peaks to specific transitions in the band structure and also to specific regions of the samples. This revealed the nature of various, previously measured interlayer exciton peaks. We further show that, due to specific optical selection rules, g-factors in vdW heterostructures are strongly spin and stacking dependent. The presented approach can potentially be applied to a wide variety of semiconductors. [1] Recent measurements of individual bands g-factors in 1L TMDs additionally corroborate the validity and accuracy of our method. [2]

## THEORY

Band structure Hamiltonian in external magnetic field:

$$H(\mathbf{B}) = H^0 + \mu_B \mathbf{B} \cdot (\mathbf{L} + \frac{g_0}{2} \boldsymbol{\Sigma}) + \frac{e_0^2}{8m_0} (\mathbf{B} \times \mathbf{r})^2$$

$\mu_B = \hbar c_0 / 2m_0$  – Bohr magneton,  $\mathbf{L} = (\mathbf{r} \times \mathbf{p}) / \hbar$  – angular momentum operator,  $\boldsymbol{\Sigma} = (\Sigma^x, \Sigma^y, \Sigma^z)$  – Pauli matrices vector.

Bloch state energy in perpendicular magnetic field:

$$\epsilon_{n\mathbf{k}}(\mathbf{B}) = \epsilon_{n\mathbf{k}}^0 + \mu_B B (L_{n\mathbf{k}} + \Sigma_{n\mathbf{k}}) + H_{n\mathbf{k}}^Q$$

where matrix elements  $L_{n\mathbf{k}} = \langle n\mathbf{k} | L^z | n\mathbf{k} \rangle$ ,  $\Sigma_{n\mathbf{k}} = \langle n\mathbf{k} | \Sigma^z | n\mathbf{k} \rangle$  and

$$H_{n\mathbf{k}}^Q = e_0^2 B^2 / 8m_0 \langle n\mathbf{k} | (r^x)^2 + (r^y)^2 | n\mathbf{k} \rangle$$

Effective g-factor of a Bloch state:

$$g_{n\mathbf{k}} = L_{n\mathbf{k}} + \Sigma_{n\mathbf{k}}$$

Orbital angular momentum matrix elements:

$$L_{n\mathbf{k}} = \frac{1}{im_0} \sum_{m=1, m \neq n}^N \frac{p_{nmk}^x p_{mnk}^y - p_{nmk}^y p_{mnk}^x}{\epsilon_{n\mathbf{k}} - \epsilon_{m\mathbf{k}}}$$

where  $r_{nmk}^\alpha = \frac{\hbar}{im_0} \frac{p_{nmk}^\alpha}{\epsilon_{n\mathbf{k}} - \epsilon_{m\mathbf{k}}}$ ,  $\epsilon_{n\mathbf{k}} \neq \epsilon_{m\mathbf{k}}$ . The coupling of the SOC term to the vector potential is taken into account by replacing the momentum operator by  $\boldsymbol{\pi} = \mathbf{p} + \frac{\hbar}{4m_0 c^2} \boldsymbol{\Sigma} \times \nabla$

Exciton energy in perpendicular magnetic field:

$$E_{\mathbf{k}}(B) = \epsilon_{c\mathbf{k}}(B) - \epsilon_{v\mathbf{k}}(B) - E_{\mathbf{k}}^{\text{bind}} = E_{\mathbf{k}}^0 + E_{\mathbf{k}}^L(B) + E_{\mathbf{k}}^Q(B)$$

Zeeman **shift** of exciton energy:

$$E_{\mathbf{k}}^L(B) = (g_{c\mathbf{k}} - g_{v\mathbf{k}}) \mu_B B = g_{\mathbf{k}} \mu_B B$$

where  $g_{\mathbf{k}}$  is **intravalley g-factor**.

Zeeman (valley) **splitting** of exciton energies:

$$\Delta E_Z = E_{\sigma+}(B) - E_{\sigma-}(B) = g \mu_B B$$

where  $g$  is **intervalley g-factor**.

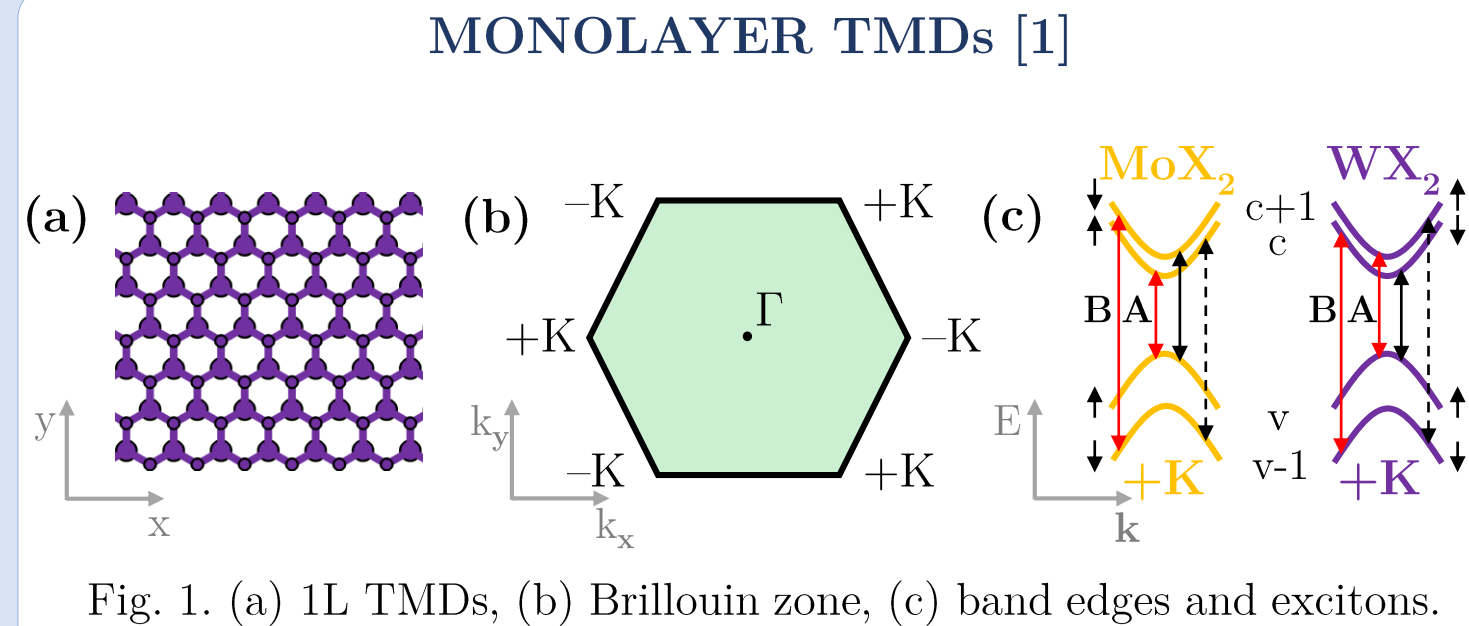


Fig. 1. (a) 1L TMDs, (b) Brillouin zone, (c) band edges and excitons.

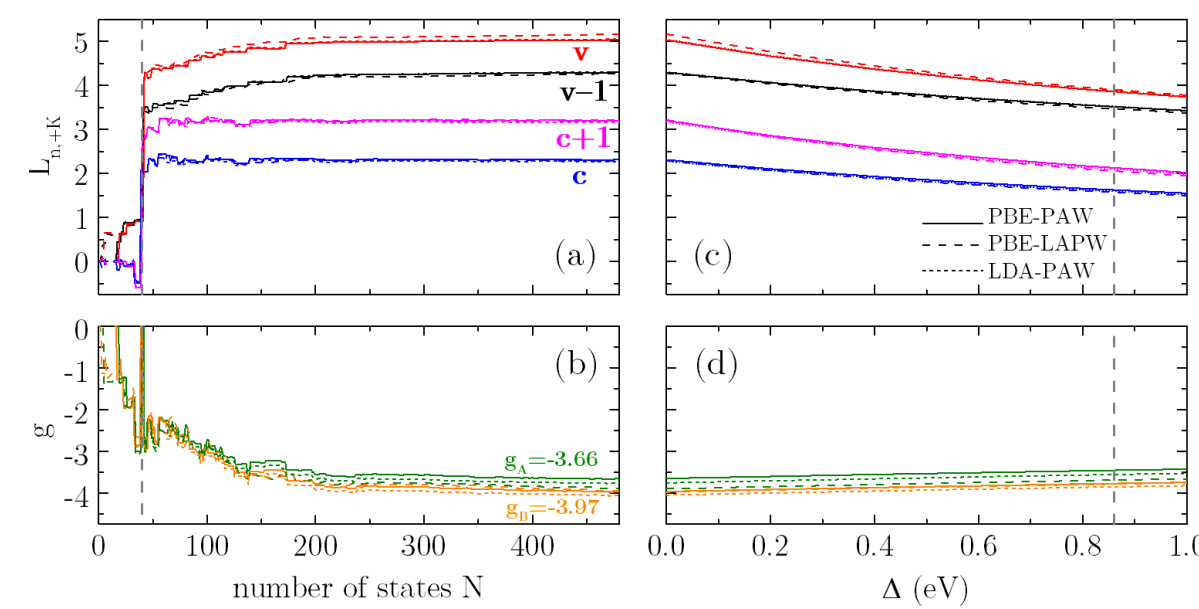


Fig. 2. WS<sub>2</sub>: (a) orbital momenta and (b) g-factors vs number of states, and vs scissor correction (c), (d) for LDA/PBE and PAW/LAPW.

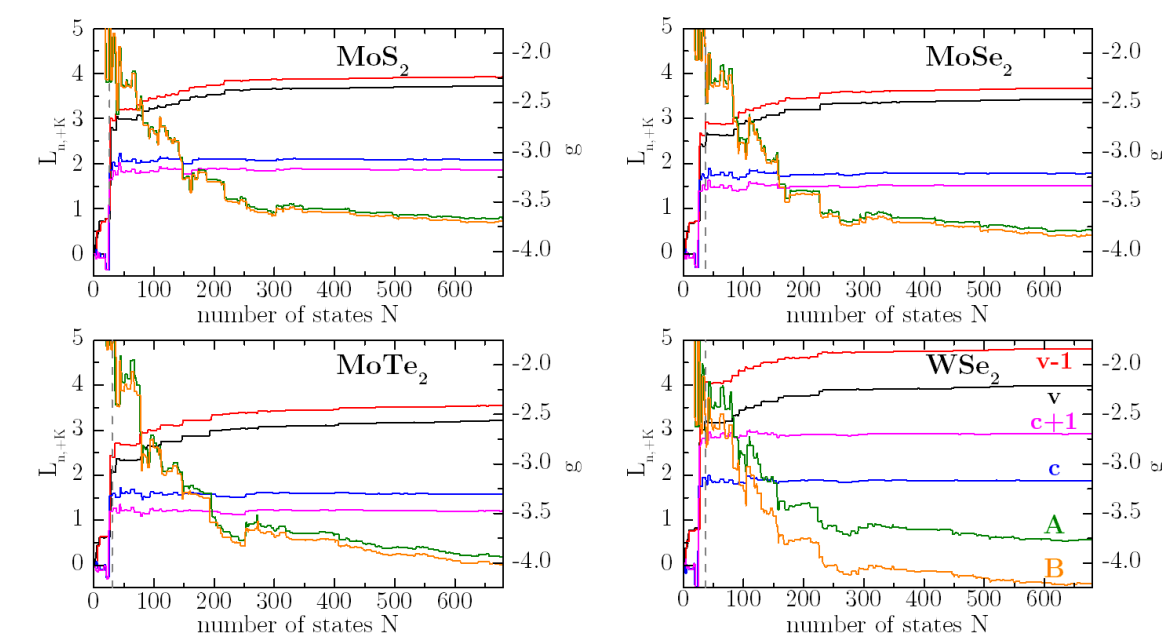


Fig. 3. Convergence of orbital momenta and g-factors for 1L TMDs.

	MoS <sub>2</sub>	MoSe <sub>2</sub>	MoTe <sub>2</sub>	WS <sub>2</sub>	WSe <sub>2</sub>
$g_{\text{B}}^{\text{L}}$	-3.68	-3.82	-3.96	-3.66	-3.80
$g_{\text{B}}^{\text{L}}(\text{exp.})$	-1.7 <sup>a</sup> -1.8 <sup>a</sup> -2.9 <sup>b</sup> -3.0 <sup>b</sup> -3.6 <sup>b</sup> -3.8 <sup>b</sup> -4.0 <sup>c</sup> -4.2 <sup>c</sup> -4.6 <sup>d</sup>	-3.8 <sup>a</sup> -4.0 <sup>a</sup> -4.1 <sup>b</sup> -4.2 <sup>b</sup> -4.3 <sup>b</sup> -4.4 <sup>d</sup>	-4.3 <sup>a</sup> -4.7 <sup>a</sup> -4.8 <sup>b</sup>	-3.7 <sup>a</sup> -3.94 <sup>a</sup> -4.0 <sup>b</sup> -4.25 <sup>c</sup> -4.35 <sup>b</sup>	-1.57 <sup>a</sup> -2.86 <sup>a</sup> -3.2 <sup>a</sup> -3.7 <sup>a</sup> -3.8 <sup>b</sup> -4.1 <sup>c</sup> -4.25 <sup>a</sup> -4.3 <sup>b</sup> -4.37 <sup>c</sup>
$g_{\text{B}}^{\text{L}}$	-3.70	-3.88	-4.02	-3.96	-4.26
$g_{\text{B}}^{\text{L}}(\text{exp.})$	-4.3 <sup>d</sup> -4.65 <sup>e</sup>	-4.2 <sup>b</sup> -3.8 <sup>b</sup>	-3.9 <sup>b</sup> -3.9 <sup>b</sup>	-3.99 <sup>a</sup> -4.9 <sup>b</sup>	-3.9 <sup>b</sup>
$L_z/L_{z+1}$	2.09/1.87	1.78/1.51	1.58/1.21	2.31/3.20	1.87/2.91
$L_{z-1}/L_z$	3.72/3.93	3.45/3.69	3.22/3.56	4.29/5.03	4.00/4.81
Intensity (Å)	28.6	21.2	13.9	42.9	33.1

Tab. 1. Calc. and exp. g-factors, orbital momenta and intensities.

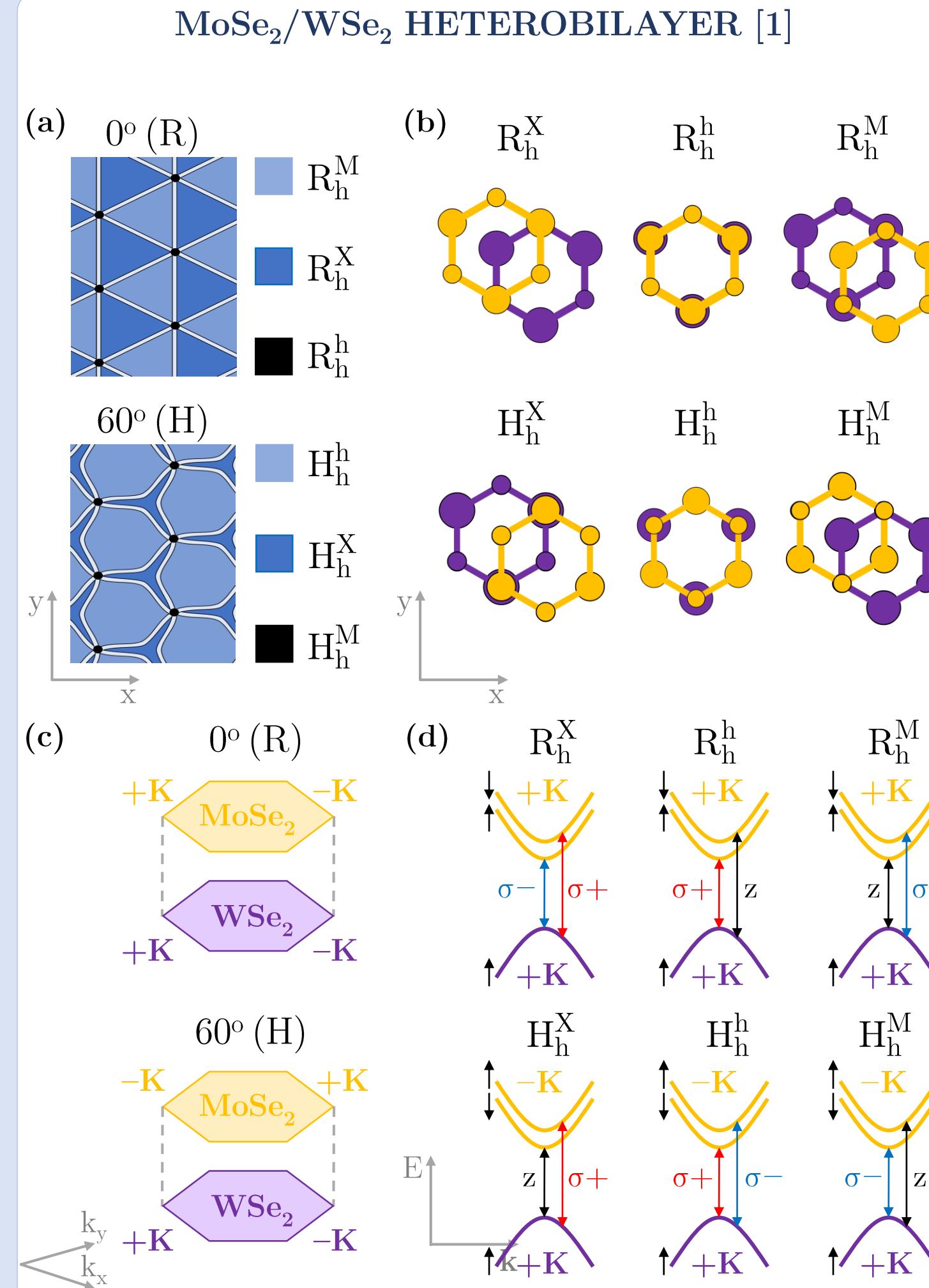


Fig. 4. MoSe<sub>2</sub>/WSe<sub>2</sub>: (a) atomic structures, (b) high symmetry stackings, (c) Brillouin zones, (d) optical selection rules for interlayer excitons.

	$R_{\text{h}}^{\text{X}}$	$R_{\text{h}}^{\text{h}}$	$R_{\text{h}}^{\text{M}}$	$H_{\text{h}}^{\text{X}}$	$H_{\text{h}}^{\text{h}}$	$H_{\text{h}}^{\text{M}}$
$g_{\text{B}}^{\text{H}}$	6.19	-10.73	-6.15	10.42	-12.60	-16.67
$g_{\text{B}}^{\text{H}}(\text{exp.})$	6.72 <sup>a</sup> 7.1 <sup>b</sup> 6.99 <sup>c</sup>	-10.6 <sup>d</sup> -8.5 <sup>b</sup>	-	-	-15.89 <sup>a</sup> -15.1 <sup>c</sup> -15.2 <sup>d</sup>	12.15 10.7 <sup>d</sup>
Transition	$v \rightarrow c$	$v \rightarrow c+1$	$v \rightarrow c$	$v \rightarrow c+1$	$v \rightarrow c+1$	$v \rightarrow c$
Intensity	0.08	0.05	0.12	$10^{-7}$	0.01	0.03
Polarization	$\sigma -$	$\sigma +$	$\sigma +$	$\sigma -$	$\sigma +$	$\sigma -$
Spin	$\uparrow \uparrow$	$\uparrow \downarrow$	$\uparrow \uparrow$	$\uparrow \downarrow$	$\uparrow \downarrow$	$\uparrow \downarrow$
$\Delta \Sigma$	0	-2	0	-2	0	-2
$L_{z(+)}$	1.80	1.53	1.79	1.53	-1.53	-1.53
$L_z$	4.90	4.90	4.86	4.74	4.77	4.54
$\Delta L$	-3.10	-3.37	-3.08	-3.21	-6.30	-6.34

Tab. 2. Calc. and exp. g-factors, orbital momenta, sel. rules and intensities.

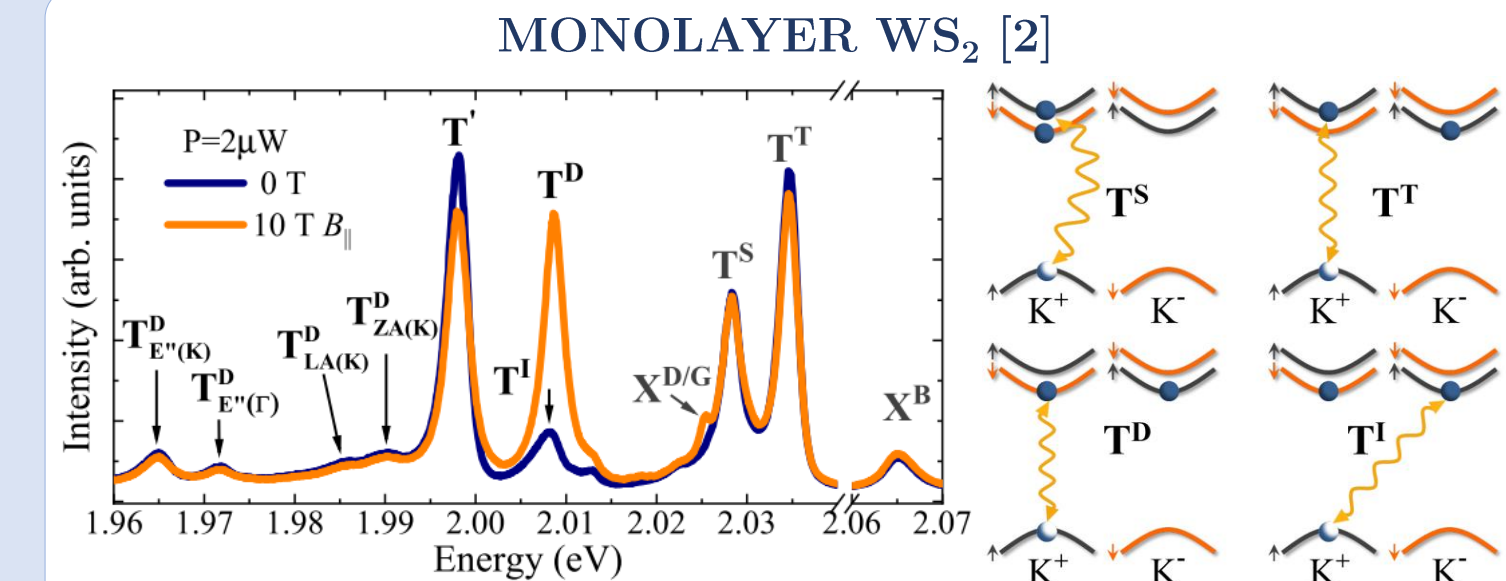


Fig. 5. 1L WS<sub>2</sub>: (a) magneto-PL spectrum, (b) K+/- trions configurations.

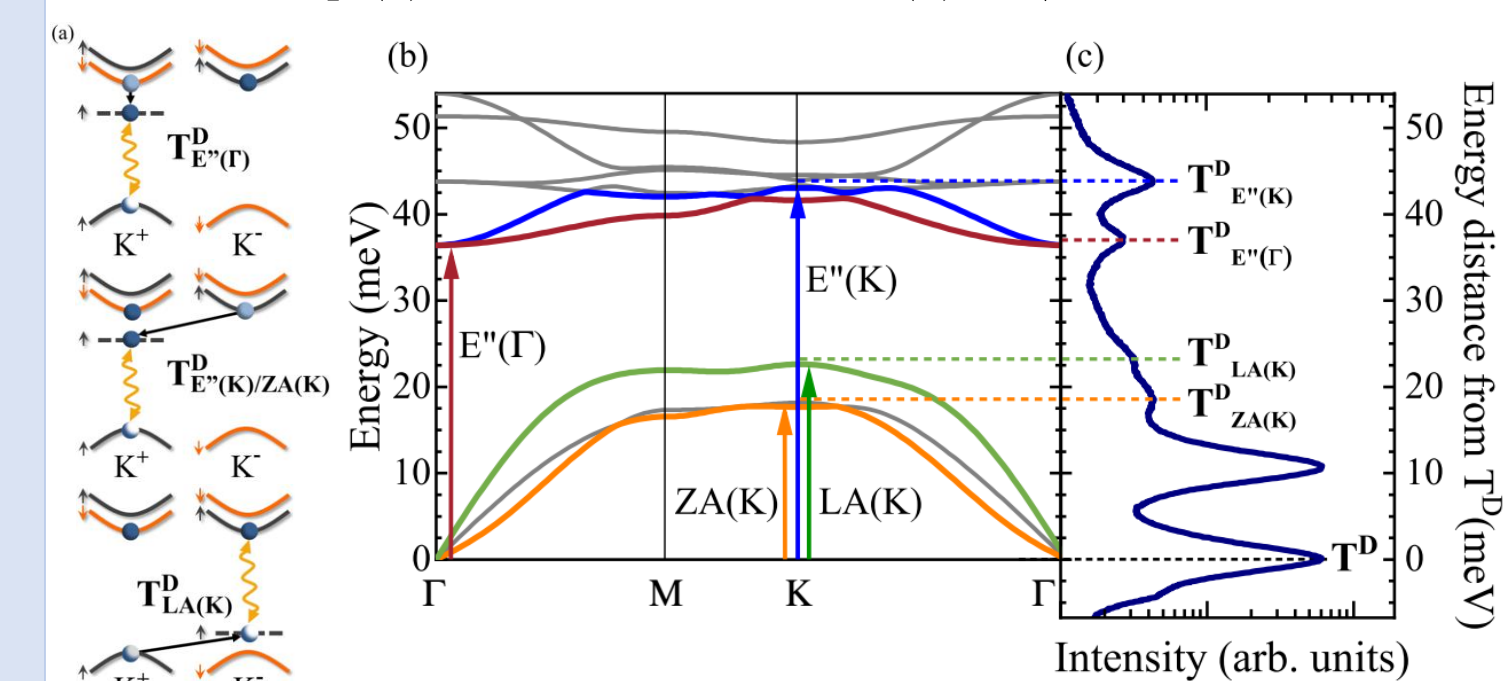


Fig. 6. Phonon replicas of TD: (a) configurations, (b) phonon dispersion, (c) magneto-PL spectrum.

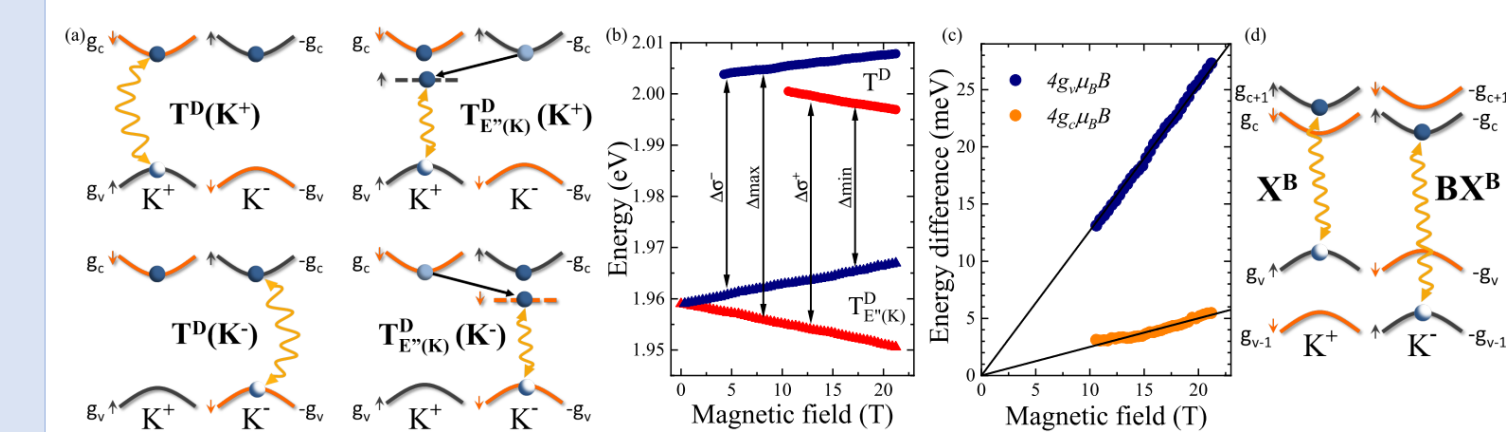


Fig. 7. Band g-factors: (a) TD & TD<sub>E(K)</sub> configurations, (b) their Zeeman shifts and (c) splittings, (d) A & B exciton configurations.

	$g_{\text{B}}^{\text{exp}}$	$g_{\text{B}}^{\text{calc}}$	$\Delta L$	$\Delta \Sigma$	$P$	$n$	$g_n^{\text{exp}}$	$g_n^{\text{calc}}$	$L_n$	$\Sigma_n$
X <sup>B</sup>	-3.5	-	-	-	-	v-1	2.81	2.79	3.79	-1
T <sup>S</sup>	-4.0	-	-	-	-	v	6.10	5.23	4.23	+1
T <sup>T</sup>	-3.9	-3.56	-1.78	0	$\sigma+$	c	0.86	0.87	1.87	-1
T <sup>I</sup>	-3.3	-	-	-	-	c+1	3.84	3.45	2.45	+1
XX <sub>1</sub> <sup>-</sup>	-4.1	-	-	-	-	-	-	-	-	-
XX <sub>2</sub> <sup>-</sup>	-4.1	-	-	-	-	-	-	-	-	-
TD	-8.9	-8.73	-2.37	-2	$\sigma+$	-	-	-	-	-
TD <sub>E(K)</sub>	-8.9	-	-	-	-	-	-	-	-	-
T <sup>I</sup>	-13.7	-	-	-	-	-	-	-	-	-
TD <sub>E(K)</sub>	-13.6	-12.20	-6.10	0	$\sigma+$	-	-	-	-	-
TD <sub>ZA(K)</sub>	-13.3	-	-	-	-	-	-	-	-	-
TD <sub>LA(K)</sub>	+13.6	+12.20	-6.10	0	$\sigma-$	-	-	-	-	-

Tab 5. Calc. & exp. band g-factors.

[2] M. Zinkiewicz, T. Woźniak, P. Kapuściński, T. Kazimierzczuk, M. Potemski, M. Molas et al, Excitonic complexes in n-doped WS<sub>2</sub> monolayer, Nano Letters, 21, 6, 2519–2525 (2021).